

# Parameterization of multi-exponential NMR relaxation in terms of logarithmic moments of an underlying relaxation time distribution

O. V. Petrov<sup>a</sup>, S. Stapf<sup>a</sup>

<sup>a</sup>Technische Universität Ilmenau, Institut für Physik, D-98684, Ilmenau, Germany

There are two common approaches to the mathematical description of multi-exponential relaxation. One is to fit an empirical function whose parameters relate to a characteristic time of relaxation and its deviation from exponential. Kohlrausch-William-Watson, Cole-Cole, Cole-Davidson and log-normal fit functions are but some of the examples. This approach requires specifying a functional form of relaxation, or in other words, *a priori* knowledge of underlying relaxation time distribution  $g(\tau)$ , which makes it difficult to compare samples with dissimilar relaxation behavior using this method. The other approach is calculation of  $g(\tau)$  by resolving an experimental relaxation function into a number of exponentials by means of the inverse Laplace transform (ILT). It provides an ultimate description of the multi-exponential relaxation and allows a direct comparison of different samples. However, since ILT of a noisy data is a numerical approximation and relies strongly on the regularization algorithm used, one may consider thus obtained  $g(\tau)$  and derived moments ambiguous.

R. Zorn [1] has proposed to characterize multi-exponential relaxation in terms of logarithmic moments (LMs) of  $g(\tau)$  which can be calculated without inversion of data by using only integral (and optionally differential) calculus. The first logarithmic moment,  $\langle \ln \tau \rangle$ , represents the geometric mean relaxation time, while the second,  $\sigma_{\ln \tau}^2$ , and the third,  $\mu_{3 \ln \tau}$ , moments relate to the width and the asymmetry of  $g(\tau)$ , respectively. The calculation of the LMs is based on the moment rules for convolution (see e.g. [2]) and requires that a relaxation function is sampled at logarithmically spaced delays and normalized to unity.

Ref. [1] contains all necessary equations for the LMs calculation and provides examples for a number of empirical fit functions used in dielectric relaxometry. Here we have applied the method to experimental data acquired under various conditions from NMR  $T_1$  relaxation and stimulated echo experiments as well as to simulated data. In particular, we tested the effects of a signal-to-noise ratio (SNR), the number of acquisition points, the onset of acquisition, and a sampling rate on the calculated LMs. It was possible to exclude a numerical differentiation from the working equations in [1], which left numerical integration and normalization of the data the only intrinsic sources of error. To normalize the data, we used the total amplitude of the best-fit triple-exponential function as a normalization factor. The same fit function was utilized to extrapolate the relaxation function to short relaxation delays if necessary.

The LMs calculated from the experimental datasets were found in a good agreement with those from ILT (UPEN) and multi-exponential analysis. The geometric mean  $\langle \ln \tau \rangle$  is robust with respect to the presence of noise, including outliers, to a moderate shift of the acquisition onset to longer delays, and to its early truncation. The variance  $\sigma_{\ln \tau}^2$  and the asymmetry parameter  $\mu_{3 \ln \tau}$  were found to be much more sensitive to those factors. For example, it was important for a good accuracy that the relaxation decay be sampled starting above a 90% level of the relaxation function, otherwise extrapolation to short delays was necessary.

This work is one of objectives of EU programme “Improving diagnosis by fast field-cycling MRI” (IDentIFY). It is intended for developing a reliable algorithm for parameterization of non-exponential relaxation processes, in order to establish a common basis for characterizing relaxation in bio-tissues and possibly to increase the contrast in  $T_1$  dispersion measurements.

## References

- [1] R. Zorn, J. Chem. Phys. 116 (2002) 3204-3209.
- [2] C. A. Laury-Micoulaud, Astron. & Astrophys. 51 (1976) 343-346.